

1

# THE QUARK-QUARK CORRELATOR:

## THEORY AND PHENOMENOLOGY

E. DI SALVO

DIPARTIMENTO DI FISICA UNIV. GENOVA

INFN - SEZ. DI GENOVA

COMO, 7-10 SEPTEMBER 2005

### SUMMARY

1. THREE PROBLEMS  
IN HIGH ENERGY REACTIONS
2. THE CORRELATOR
3. SPLITTING INTO A T-EVEN  
AND A T-ODD PART
4. EQUATIONS OF MOTION
5. CONCLUSIONS

# 1. THREE PROBLEMS IN HIGH ENERGY REACTIONS

- DETERMINING TRANSVERSITY
- AZIMUTHAL ASYMMETRIES
- FUNCTION  $g_2$

## A) TRANSVERSITY

$$h_1(x) = \int \frac{d\lambda}{2\pi} e^{i\lambda x} \langle P, S | \bar{\Psi}(0) \gamma_5 \gamma^+ \gamma_\perp \cdot S \mathcal{L}(\lambda) \Psi(\lambda) | P, S \rangle$$

$$\mathcal{L}(\lambda) = \mathbb{P} \exp \left\{ -i g \int_0^\lambda A^+(x') dx' \right\}$$

AXIAL GAUGE  $\Rightarrow \mathcal{L}(\lambda) = 1 \quad h_1(x) = q^+(x) - q^-(x)$

$h_1(x)$  CHIRAL ODD  $\Rightarrow$  ~~DIS~~

"DIRECT" METHODS:

$\left\{ \begin{array}{l} p^\uparrow p^\uparrow \\ p^\uparrow p \end{array} \right.$	$\begin{array}{l} \nearrow \mu^+ \mu^- X \\ \searrow J/\psi X \end{array}$	$G \propto h_1 \otimes \bar{h}_1$
	$\rightarrow \Xi^\uparrow X$	$G \propto h_1^p \otimes h_1^{\Xi}$

"INDIRECT" METHODS:

$\left\{ \begin{array}{l} p^\uparrow \\ p^\uparrow \end{array} \right.$	$\begin{array}{l} \nearrow l\pi X \\ \searrow l\pi\pi X \end{array}$	$G \propto h_{1T} \otimes H_1^\perp$
		$G \propto h_{1T} \otimes (I.F.)$

$h_1(x) = \int d^2 p_\perp h_{1T}(x, \vec{p}_\perp^2) \quad H_1^\perp, (I.F.) : T\text{-ODD}$

## B) AZIMUTHAL ASYMMETRIES

- $l p^\uparrow \rightarrow l \pi X$  HERMES, COMPASS, CLAS
- $p p(\pi) \rightarrow \mu^+ \mu^- X$  ( $\cos 2\varphi$  ASYMM.) NA10
- $p N \rightarrow Y^\uparrow X$  ( $Y = \text{HYPERON}$ )  
 $p p^\uparrow \rightarrow \pi^\circ X$  RHIC

SUCH ASYMMETRIES CANNOT  
BE EXPLAINED BY PQCD

### POSSIBLE EXPLANATIONS:

- QUARK-QUARK-GLUON CORRELATIONS

J. Qiu & G. Sterman: P.R.L. 67 (1991) 2264; Nucl. Phys. B 378 (1992) 52

- TRANSVERSE MOMENTUM  $\rightarrow$  T-ODD FUNCTIONS

D. W. Sivers: P.R. D 41 (1990) 83; P.R. D 43 (1991) 261

D. Boer, R. Jakob and P.J. Mulders: Nucl. Phys. B 564 (2000) 471

M. Anselmino et al.: P.R. D 71 (2005) 074006

J. C. Collins: Ph. Lett. B 536 (2002) 43

DISCRETE AGREEMENT WITH DRELL-YAN

BUT NOT REPRODUCED  $M_{\text{eff.}}$  DEPENDENCE

## C) FUNCTION $g_2(x)$

A PUZZLING OBJECT

WANDZURA - WILCZEK APPROXIMATION:

$$g_1(x) + g_2(x) = \int_x^1 \frac{dy}{y} g_1(y)$$

OPERATOR PRODUCT EXPANSION  $\Rightarrow$

$$g_2(x) = g_2^{(2)}(x) + g_2^{(3)}(x) \quad (2), (3) \rightarrow \text{TWIST}$$

$$\int_0^1 dx x^{n-1} \left\{ g_1(x) + \frac{n}{n-1} g_2^{(2)}(x) \right\} = 0$$

$n \geq 3$  :

$$\int_0^1 dx x^{n-1} \left\{ g_1(x) + \frac{n}{n-1} g_2^{(3)}(x) \right\} = \frac{1}{2} \sum_i S_i d_n^i E_{3,i}^n$$

A PUZZLE:

$$\text{OPE} \Rightarrow g_2^{(2)}(x) + g_2(x) \neq 0$$

INSTEAD Jaffe & Ji: P.R.L. 67 (1991) 552; M.P.B. 375  
(1992) 527

$$\Rightarrow g_2(x) + g_1(x) = (\text{TWIST } 3)$$

## 2. DEFINITION OF THE CORRELATOR

REACTIONS AT HIGH MOMENTUM TRANSFER:

$$d\sigma \propto (\text{Soft}) \otimes (\text{Hard})$$

FACTORIZATION THEOREMS:

J. C. Collins: Phys. Rev. D 57 (1998) 3051

J. C. Collins & D. E. Soper: Nucl. Phys. B 193 (1981) 381;

Nucl. Phys. B 197 (1982) 446

X. Ji, J. P. Ma & F. Yuan: Phys. Lett. B 597 (2004) 299;

Phys. Rev. D 71 (2005) 034005

"SOFT" FUNCTIONS ENCODED IN

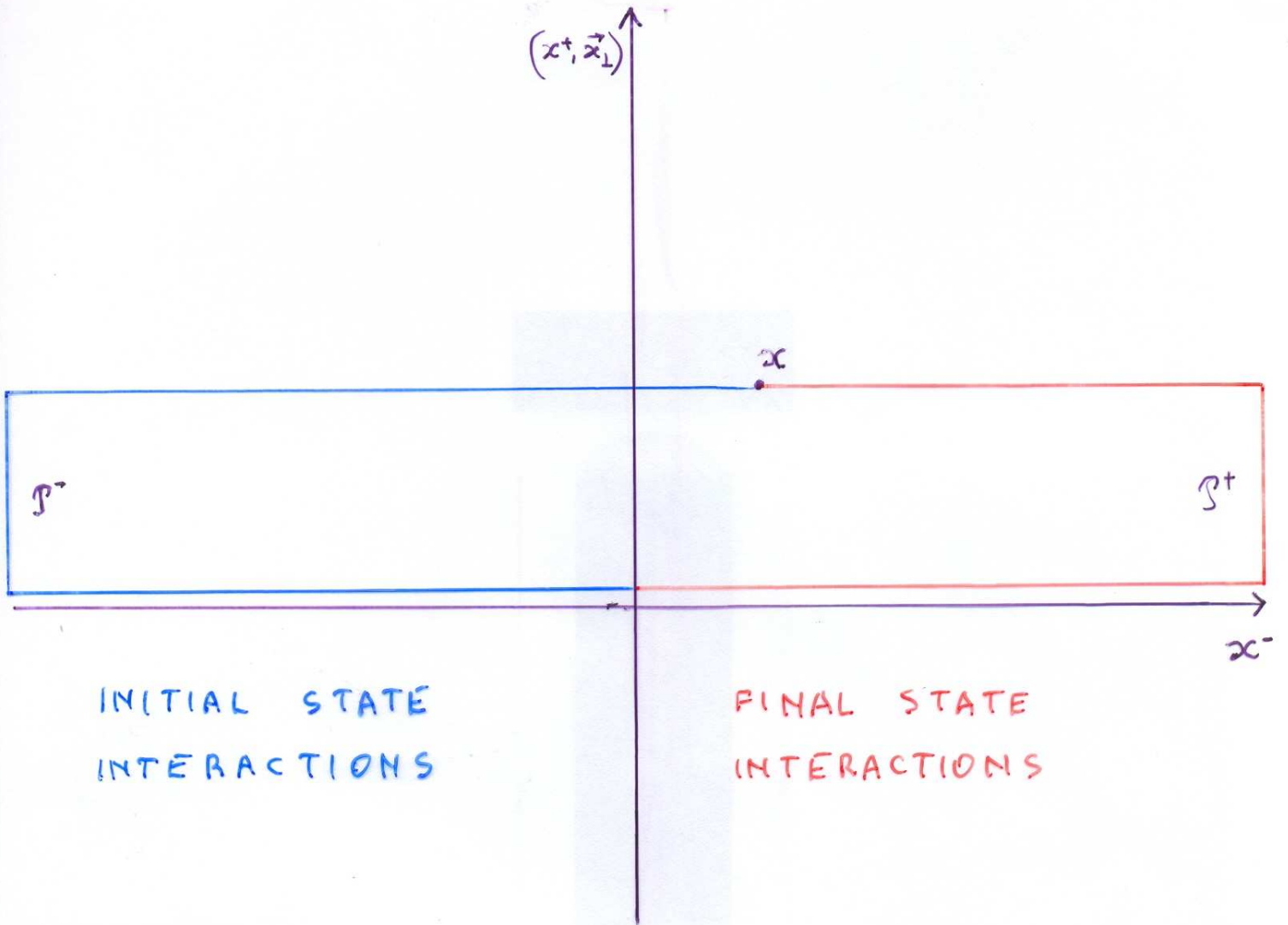
CORRELATOR:

$$\Phi_{ij} = \int \frac{d^4x}{(2\pi)^4} e^{ipx} \langle P, S | \bar{\psi}_j(0) \mathcal{L}(x) \psi_i(x) | P, S \rangle$$

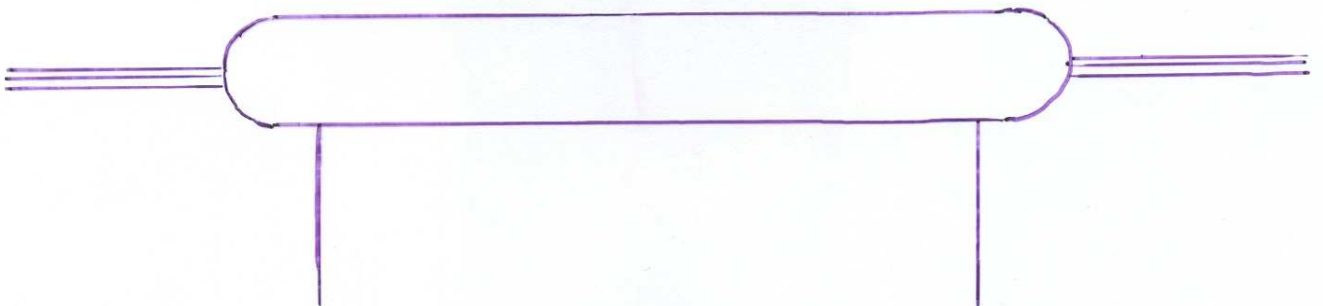
$$\mathcal{L}(x) = \mathbb{P} \exp \left\{ -ig \int_0^x \frac{dz}{(P^+)} A_a^\mu(z) \lambda^a dz_\mu \right\}$$

R. J. Ralston & D. E. Soper: Nucl. Phys. B 152 (1979) 109

P. J. Mulders & R. D. Tangerman: Nucl. Phys. B 461 (1996) 197



$$\Phi_{ij} = \int \frac{d^4x}{(2\pi)^4} e^{ip \cdot x} \langle P, S | \bar{\psi}_j(0) \mathcal{L}(x) \psi_i(x) | P, S \rangle$$



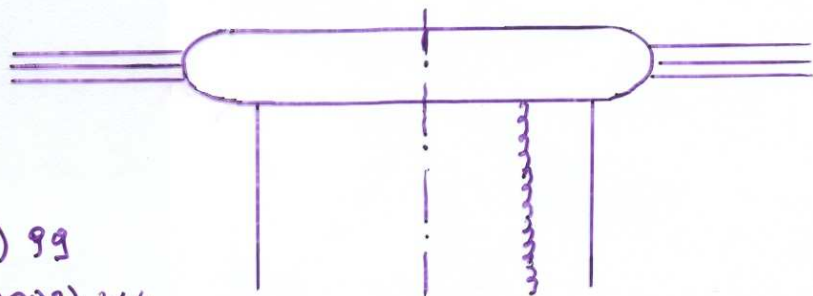
$\Phi$ : GREEN FUNCTION : 2 NUCLEON LEGS + 2 QUARK LEGS  
+ 0, 1, 2 ...  $n$  ... GLUON LEGS

# PROPERTIES OF $\phi$

- HERMITICITY OF  $\gamma_0 \phi$
- P-INVARIANCE
- BUT T-EVEN AND T-ODD FUNCTIONS

T-ODD FUNCTIONS IN STRONG INTERACTIONS:

INTERFERENCE



S. Brodsky et al.: P.L. B 530 (2002) 99  
Nuc. Phys. B 642 (2002) 344

DIRAC ALGEBRA: 16 ELEMENTS



32 FUNCTIONS IN ALL:

K. Goebel, A. Metz & M. Schlegel: hep-ph/0504130

→ Ph. Lett. B 618 (2005) 90

# DISTRIBUTION FUNCTIONS

PROJECTIONS ONTO DIRAC

COMPONENTS :

$$\Phi^\Gamma = \frac{1}{2} \int dp^- \text{tr} (\Phi \Gamma) \quad \Gamma \in \text{DIRAC ALG.}$$

$$\Phi^{\gamma^+} \supset f_1(x, \vec{p}_\perp^2) \quad \text{UNPOLARIZED DENSITY}$$

$$\Phi^{\gamma_s \gamma^+} = g_{1L}(x, \vec{p}_\perp^2) \quad \text{LONG. POL. DENSITY}$$

$$\Phi^{\gamma_s \gamma^+ \gamma^i} \supset h_{1T}(x, \vec{p}_\perp^2) \quad \text{TRANSVERSITY}$$



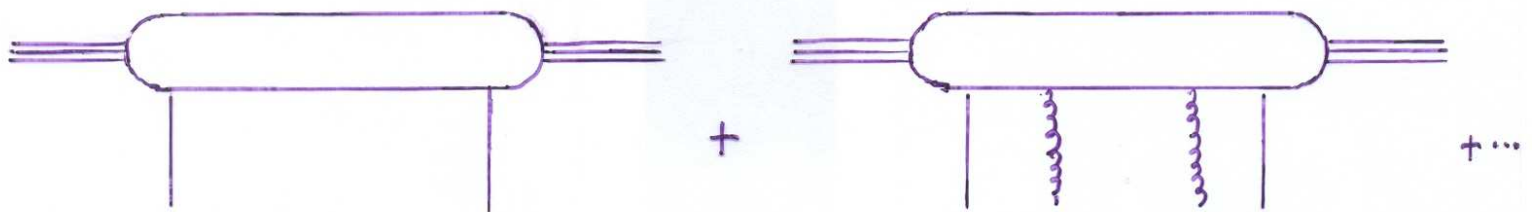
### 3. SPLITTING OF $\phi$ INTO A T-EVEN AND A T-ODD CONTRIBUTION

$$L_E(x) = \mathbb{P} \cos \left\{ g \int_{o(\mathcal{S}^+)}^x A_a^\mu(z) \lambda^a dz_\mu \right\}$$

$$L_O(x) = -i \mathbb{P} \sin \left\{ g \int_{o(\mathcal{S}^+)}^x A_a^\mu(z) \lambda^a dz_\mu \right\}$$

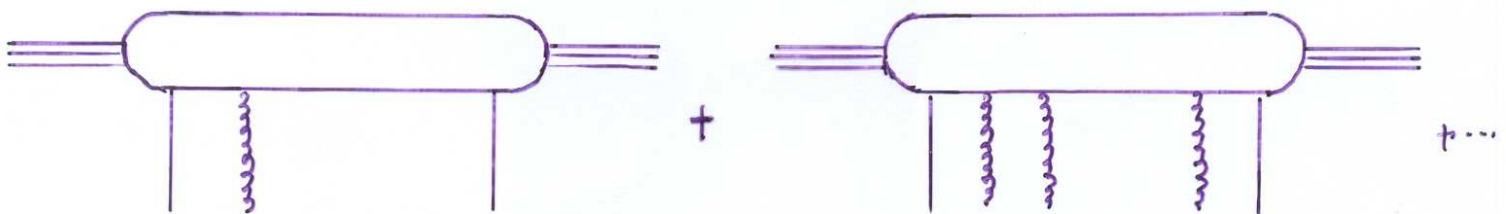
$$L_E(x) \rightarrow \phi_E = \phi_E^{(0)} + g^2 \phi_E^{(2)} + g^4 \phi_E^{(4)} + \dots : \text{T-EVEN}$$

$\mathcal{S}^+ + \mathcal{S}^-$



$$L_O(x) \rightarrow \phi_O = g \phi_O^{(1)} + g^3 \phi_O^{(3)} + g^5 \phi_O^{(5)} + \dots : \text{T-ODD}$$

$\mathcal{S}^+ - \mathcal{S}^-$



## "SOFT" FUNCTIONS

- T-EVEN FUNCTIONS:  $f_E = \frac{1}{2} \text{tr}(\phi_E \Gamma) = \phi_E^\Gamma$

IN PARTICULAR:

$$f_2(x, \vec{p}_\perp^2) = \phi_E^{\gamma^+}$$

$$g_{1L}(x, \vec{p}_\perp^2) = \phi_E^{\gamma^+ \gamma^+}$$

$$h_{1T}(x, \vec{p}_\perp^2) = \phi_E^{\gamma^+ \gamma^+ \gamma_\perp \cdot S}$$

- T-ODD FUNCTIONS:  $f_0 = \frac{1}{2} \text{tr}(\phi_0 \Gamma) = \phi_0^\Gamma$

IN PARTICULAR:

- $f_{1T}^\perp(x, \vec{p}_\perp^2) \varepsilon^{\alpha\beta\gamma} \frac{p_\perp^\alpha}{\mu} \frac{p_\perp^\beta}{|\vec{p}_\perp|} S^\gamma = \phi_0^{\gamma^+}$

UNPOLARIZED QUARK DENSITY IN A

TRANSVERSELY POLARIZED NUCLEON (SIVERS FUNC.)

- $h_1^\perp(x, \vec{p}_\perp^2) = \phi_0^{\gamma^+ \gamma^+ \gamma_\perp \cdot z_\perp} \quad z_\perp^\alpha = \varepsilon^{\alpha\beta\gamma} \frac{p_\perp^\alpha}{\mu} \frac{p_\perp^\beta}{|\vec{p}_\perp|}$

QUARK TRANSVERSITY IN AN UNPOLARIZED

NUCLEON ("COLLINS" FUNCTION)

## 4. EQUATIONS OF MOTION

## POLITZER'S THEOREM

(H. D. Politzer: Nucl. Phys. B 172 (1980) 349)

↓

$$\langle P, S | \bar{\psi}_i(0) \not{L}(x) (\not{x} - m_q)_{ij} \psi_j(x) | P, S \rangle = 0$$

↓

SETTING

$$\phi_E = \phi_E^{(0)} + g^2 \phi_E^{(2)} + \dots, \quad \phi_0 = g \phi_0^{(1)} + \dots,$$

WE GET

$$(\not{p} - m_q) \phi_E^{(0)}(p) = 0,$$

$$(\not{p} - m_q) \phi_0^{(1)}(p) = \Psi_0^{(1)}(p),$$

$$(\not{p} - m_q) \phi_E^{(2)}(p) = \Psi_E^{(2)}(p),$$

WHERE, IN AN AXIAL GAUGE,

$$[\Psi_0^{(1)}(p)]_{ij} = \int \frac{d^4 x}{(2\pi)^4} e^{ipx} \langle P, S | \bar{\psi}_i(0) [A_\infty]_{ik} \psi_k(x) | P, S \rangle,$$

$$[\Psi_E^{(2)}(p)]_{ij} = g \frac{1}{\not{p} - m_q} \int \frac{d^4 x}{(2\pi)^4} e^{ipx} \langle P, S | \bar{\psi}_i(0) P\{A_\infty^2\} \psi_j(x) | P, S \rangle$$

$$A_{\infty\mu} = A_{\perp\mu}(x^- = \infty, x^+, \vec{x}_\perp)$$

## A) ZERO ORDER APPROXIMATION

$$(\not{p} - m_q) \phi_E^{(0)} = 0 \quad + \quad \text{HERMITICITY OF } \gamma_0 \phi_E$$

↓

$$\phi_E^{(0)} = \frac{\not{p} + m_q}{2} \left( f_1^{(0)} + g_{1L}^{(0)} \not{S}_{11} \gamma_5 + h_{1T}^{(0)} \not{S}_{1L} \gamma_5 \right) \delta(p^2 - m_q^2)$$

ONLY 3 INDEPENDENT FUNCTIONS

↓

RELATIONS AMONG DIFFERENT DIRAC COMPONENTS :

### TWIST 2

$$g_{1T}^{(0)} = h_{1T}^{(0)}$$

$g_{1T}$  : LONGIT. POL. QUARK DENSITY IN A  
TRANSV. POL. NUCLEON

IMPORTANT FOR DETERMINING TRANSVERSITY

$$h_{1L}^{(0)} = g_{1L}^{(0)}$$

$h_{1L}$  : TRANSV. POL. QUARK DENSITY IN A  
LONGIT. POL. NUCLEON

### TWIST 3

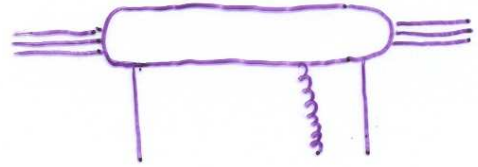
$$g_T^{(0)} = m_q h_1^{(0)}$$

$$g_T = \int \phi_E \gamma_5 \gamma_1 \cdot S \quad d^2 p_{\perp} = g_1 + g_2$$

## B) FIRST ORDER (T-ODD) CORRECTION

$$(\not{p} - m_q) \phi_0^{(1)} = \Psi_0^{(1)}$$

$$\Psi_0^{(1)} \rightarrow$$

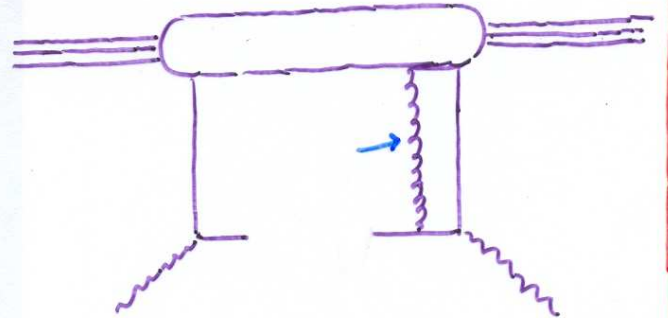
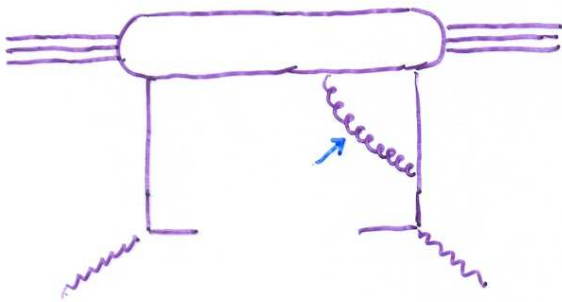


q-q-g CORRELATOR

$$\phi_0^{(1)} = (\not{p} + m_q) \Psi_0^{(1)} \mathcal{P} \frac{1}{p^2 - m_q^2} \Rightarrow$$

$$\phi_0^{(1)} \propto \frac{|\vec{p}_1^+|}{p^+}$$

EFFECTIVE T-ODD FUNCTIONS:



SOFT GLUON APPROXIMATION ( $|\vec{p}_1^+| \ll Q$ ):

J. C. Collins: P. R. D 57 (1998) 3051; P. L. B 356 (2002) 43

T-ODD FUNCTIONS  $\propto Q^{-1}$

E. G.,

SINGLY POLARIZED SIDIS ASYMM.:  $\underline{h_{1T}} \otimes \underline{H_1^\perp} \propto Q^{-1}$

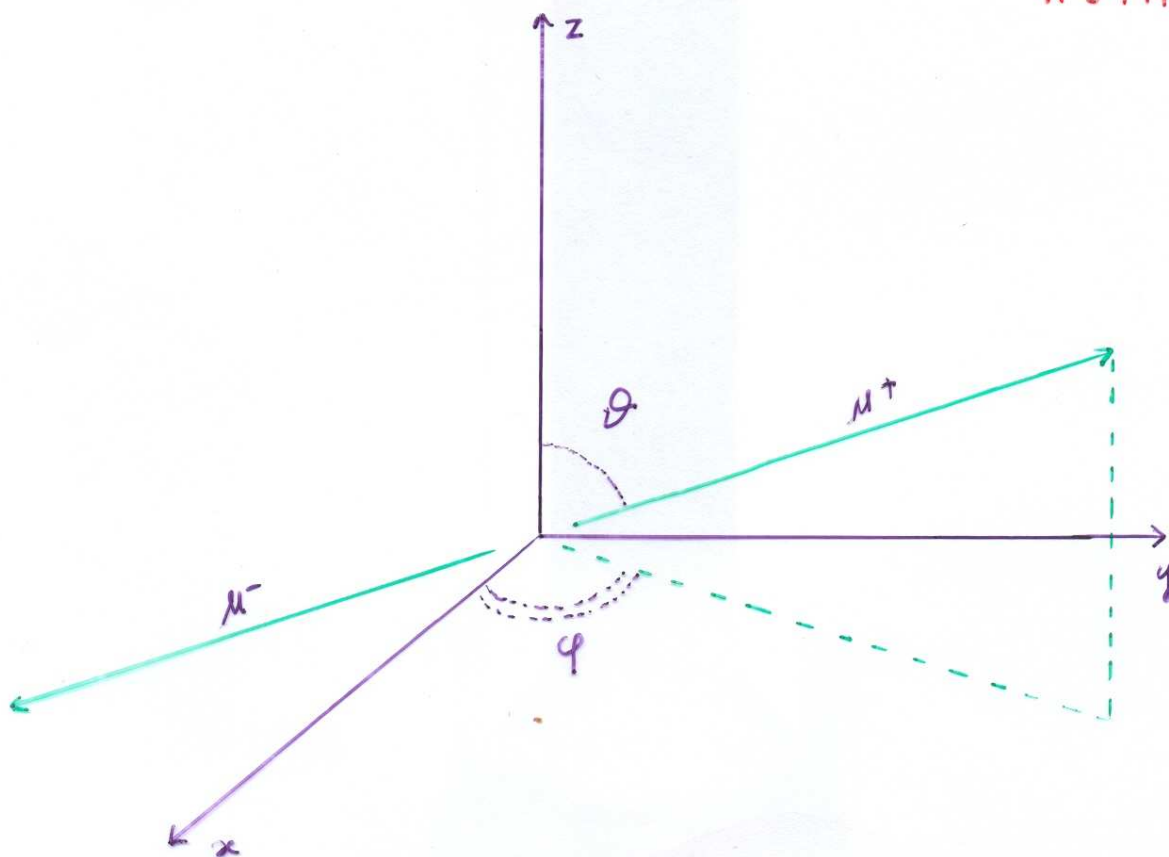
UNPOLARIZED DRELL-YAN ASYMM.:  $\underline{h_1^\perp} \otimes \underline{\bar{h}_1^\perp} \propto Q^{-2}$

## UNPOLARIZED DRELL-YAN

$$p\pi \rightarrow \mu^+ \mu^- X$$

$$\frac{1}{\sigma} \frac{d\sigma}{d\Omega} = \frac{3}{4\pi} \frac{1}{\lambda+3} \left( 1 + \lambda \cos^2 \vartheta + \mu \sin 2\vartheta \cos \varphi + \frac{1}{2} \nu \sin^2 \vartheta \cos 2\varphi \right)$$

AZIMUTHAL  
ASYMMETRY



DATA FROM NA 10 COLLAB.:

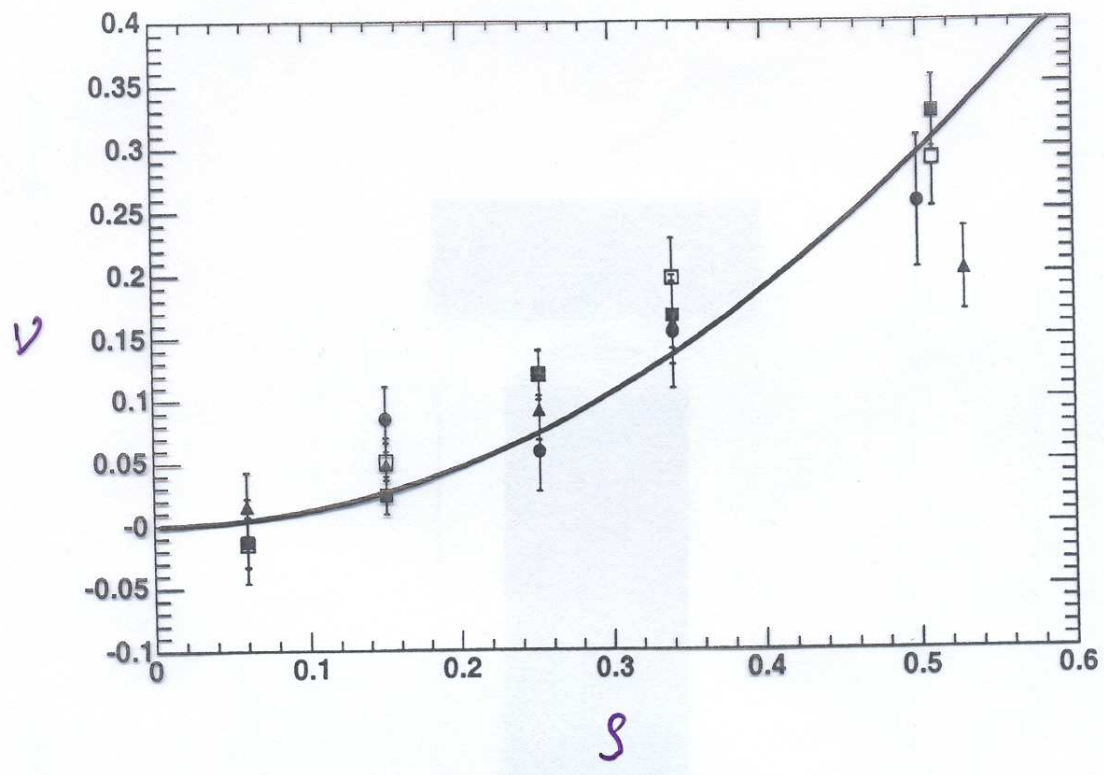
S. Falciano et al.: Z. Phys. C 31 (1986) 513

M. Guanziriodi et al.: Z. Phys. C 37 (1988) 545

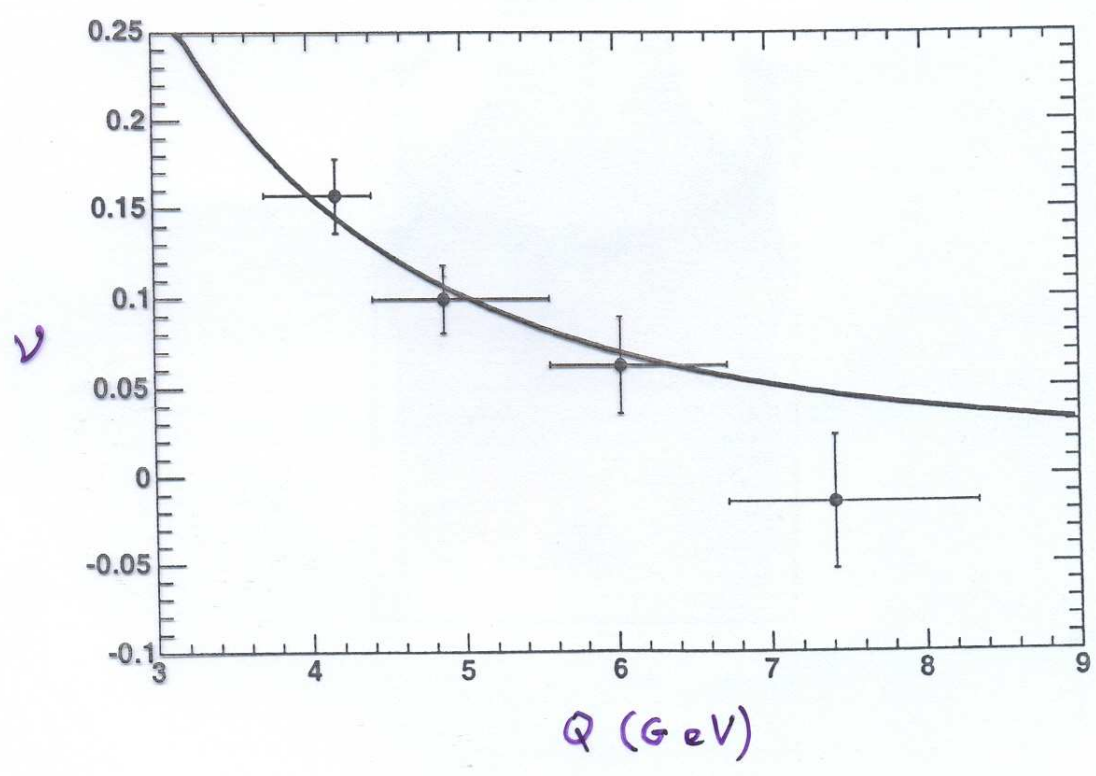
RESULT EOM:

$$\nu = A \frac{q_{\perp}^2}{Q^2}$$

$q_{\perp}$  = T. M. OF  $(\mu^+ \mu^-)$



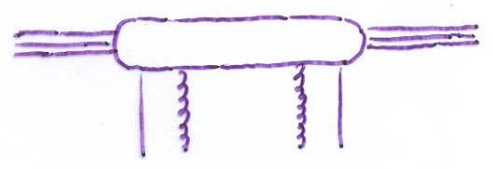
$$\beta = \frac{q^2}{Q}$$



# C) SECOND ORDER (T-EVEN) CORRECTION

$$(\not{p} - m_q) \Phi_E^{(2)} = \Psi_E^{(2)}$$

$$\Psi_E^{(2)} \rightarrow$$



qqqq CORRELATOR

$$\Phi_E^{(2)} = (\not{p} + m_q) \Psi_E^{(2)} \int \frac{1}{p^2 - m_q^2} \Rightarrow$$

$$\Phi_E^{(2)} \propto \frac{|\vec{P}_1|^2}{p^{+2}}$$

TO SUMMARIZE,  $\Phi_E$  READS AS

$$\Phi_E = \frac{p^+}{2} (\not{p} + m_q) (f_1 + g_{1L} \gamma_5 S_{11}^i + h_{1T} \gamma_5 S_{1\perp}^i) \delta(p^2 - m_q^2) + g^2 (\not{p} + m_q) \Psi_E^{(2)} \int \frac{1}{p^2 - m_q^2} + O\left(g^4 \frac{|\vec{P}_1|^4}{p^{+4}}\right)$$



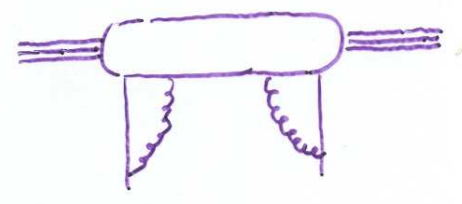
$$h_{1T} = g_{1T} + O\left(g^2 \frac{|\vec{P}_1|^2}{p^{+2}}\right) \rightarrow O\left(g^2 \frac{|\vec{P}_1|^2}{Q^2}\right)$$

$$g_T = g_1 + g_2 = m_q h_{1\perp}^{(0)} + \left[ (\not{p} + m_q) \frac{1}{2} \text{tr}(\Psi_E^{(2)} \gamma_5 \gamma_\perp \cdot S) \int \frac{1}{p^2 - m_q^2} \right] + O\left(g^4 \frac{|\vec{P}_1|^4}{p^{+4}}\right)$$

$O\left(g^2 \frac{|\vec{P}_1|^2}{Q^2}\right) \rightarrow$  PREDICTION

$\Rightarrow$  LIMITS ON OPE FOR  $g_2$

"EFFECTIVE" FUNCTIONS:  $\Phi_E \rightarrow \bar{\Phi}_E$  (E.G.)



$$g_T^{\text{eff}} = \frac{1}{2} \text{tr}(\bar{\Phi}_E \gamma_5 \gamma_\perp \cdot S) \text{ IS JUST TWIST 3}$$

$\downarrow$   
JAFFE - JI



## 5. CONCLUSIONS

### 1. APPROXIMATE EQUALITIES BETWEEN

CHIRAL ODD AND CHIRAL EVEN FUNCTIONS:

$$h_{1T}(x, \vec{p}_\perp^2) = g_{1T}(x, \vec{p}_\perp^2) + O\left(g^2 \frac{|\vec{p}_\perp^2|}{Q^2}\right) \rightarrow \text{TRANSVERSITY}$$

$$h_{1L}(x, \vec{p}_\perp^2) = g_{1L}(x, \vec{p}_\perp^2) + O\left(g^2 \frac{|\vec{p}_\perp^2|}{Q^2}\right)$$

### 2. AZIMUTHAL ASYMMETRIES:

$$qqq \text{ CORRELATIONS} \propto Q^{-1}$$

AGREEMENT WITH UNPOL. A. A. DRELL-YAN:

$$h_1 \otimes \bar{h}_1 \propto Q^{-2}$$

EFFECTIVE T-ODD FUNCTIONS  $\rightarrow$  APPROXIMATE SOLUTIONS TO E.O.M. (SOFT GLUON APPROX.)

### 3. PREDICTIONS ABOUT $g_2(x)$ :

$$g_1(x) + g_2(x) = m_q h_1^{(1)}(x) + O\left(g^2 \frac{|\vec{p}_\perp^2|}{Q^2}\right)$$

$\Rightarrow$  LIMITS ON OPE FOR TRANSV. POL. DIS

"EFFECTIVE"  $g_1 + g_2$  JUST TWIST 3